- 1. Basic design rules for electrical machines
- 2. Design of Induction Machines
- 3. Heat transfer and cooling of electrical machines
- 4. Dynamics of electrical machines
- 5. Dynamics of DC machines
- 6. Space vector theory
- 7. Dynamics of induction machines
- 8. Dynamics of synchronous machines





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Source:

SPEED program





6. Space vector theory







6. Space vector theory

- 6.1 M.M.F. space vector definition
- 6.2 M.M.F. space vector and phase currents
- 6.3 Current, flux linkage and voltage space vectors
- 6.4 Space vector transformation
- 6.5 Influence of zero sequence current system
- 6.6 Magnetic energy



6. Space vector theory Concept of space vector ("Raumzeiger")





- Stator and rotor <u>fundamental</u> air gap field are excited by sinusoidal distributed stator and rotor current load (= "current layer" $A_s(\gamma_s), A_r(\gamma_s)$)
- Superposition of both fundamental fields yields the resulting magnetizing fundamental air gap field wave $B(\gamma_s)$
- Each sinusoidal distributed air gap field wave is described by a **space vector** in the machine's axial cross section plane



6. Space vector theory Definition of space vector ("Raumzeiger")





- Space vector length = field wave amplitude \hat{B}
- Space vector orientation = position γ_s of field wave maxima
- Space vector direction = position of north pole N.
- Use of complex coordinate frame for machine cross section \Rightarrow complex space vector!
- Alternatively the space vectors $\underline{\underline{B}}$ or $\underline{\underline{\Psi}}$ or $\underline{\underline{I}}$ are used for the magnetic fundamental air gap field wave \Rightarrow e.g.: $\underline{\underline{B}}_{\delta}$ or $\underline{\underline{\Psi}}_{h}$ or $\underline{\underline{I}}_{m}$ $\hat{\underline{I}}_{m} = \underline{\underline{\Psi}}_{h} / L_{h} = (N_{s}k_{ws} \cdot \frac{2}{\pi}\tau_{p}l_{e} / L_{h}) \cdot \underline{\underline{B}}_{\delta}$



M.M.F. space vector definition



Dynamic situation:

• **Dynamic situation:** The three phase currents are no longer of sine wave time function, but $I_{\rm U}(t)$, $I_{\rm V}(t)$, $I_{\rm W}(t)$ vary arbitrarily.

Steady state:

fixed frequency Ω , amplitude \hat{I} & phase shift currents change arbitrarily

$$I_U(t) = \hat{I} \cdot \cos(\Omega \cdot t) \qquad I_U(t)$$

$$I_V(t) = \hat{I} \cdot \cos(\Omega \cdot t - 2\pi/3) \qquad I_V(t)$$

$$I_W(t) = \hat{I} \cdot \cos(\Omega \cdot t - 4\pi/3) \qquad I_W(t)$$

- $I_W(t) = I \cdot \cos(\Omega \cdot t 4\pi / 3)$
- In many cases the three-phase winding is **star-connected**: $I_{II}(t) + I_{V}(t) + I_{W}(t) = 0$
- a) **Delta-connected** winding or b) star-connected winding with **connection of neutral point**: $I_{U}(t) + I_{V}(t) + I_{W}(t) \neq 0$,,neutral point clamped"



Three-phase sinusoidal AC current system Three-phase arbitrary current system <u>Fixed</u> frequency f = 1/TNo frequency detectable <u>Fixed</u> amplitude \hat{I} No defined amplitude Fixed phase shift 0, $2\pi/3$, $4\pi/3$ No phase-shift defined $0\frac{1}{12}\frac{1}{6}$ $i_{\rm W}(t)$ 111 1 VW $i_{\rm U}(t$ 0 $i_{\rm V}(t)$ Т



6. Space vector theory

Three-phase system: Sinusoidal symmetrical vs. arbitrary currents



Common mode current $i_0(t)$ in all three phases identical!





"Common-mode free" current system:

$$i_{US}(t) + i_{VS}(t) + i_{WS}(t) = 0$$

Original current system:

$$i_U(t)+i_V(t)+i_W(t)\neq 0$$

Example: Neutral point N clamped



$$i_0(t) = [i_U(t) + i_V(t) + i_W(t)]/3$$



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6. Space vector theory Common mode current $i_0(t)$



$$i_0(t) = [i_U(t) + i_V(t) + i_W(t)]/3$$

Three-phase sinusoidal AC current system Three-phase arbitrary current system



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Example: Common mode current $i_0(t)$







Three phase sinusoidal currents





"Steady state" operation

• Fundamental m.m.f. wave $V_{\nu=1}(x_s)$ moves with constant velocity

$$v_{syn} = \frac{2\tau_p}{T} = 2f\tau_p$$

Synchronous velocity !

Synchronous rotational speed n_{syn} in case of <u>rotating</u> field arrangement:

$$\omega_{syn} = 2\pi n_{syn} = \frac{v_{syn}}{d_{si}/2} = \frac{v_{syn}}{p\tau_p/\pi} = \frac{2\pi f}{p}$$

n_{svn}



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6. Space vector theory M.M.F. V(x_s) of three phase winding at <u>arbitrary</u> currents



Three-phase AC star-connected winding with <u>arbitrary phase currents</u> excites a MMF distribution V(x) with a <u>dominant</u> sine wave fundamental $V_1(x_s)$

Example: q = 2, $N_c = 1$, single layer winding, $I_U(t) = I$, $I_V(t) = 2I$, $I_W(t) = -3I$





<u>Note:</u> Star connection: Common mode current I_0 is ZERO!

$$3 \cdot I_0(t) = I_U(t) + I_V(t) + I_W(t) = I + 2I - 3I = 0$$

Fourier m.m.f. fundamental:

$$V_1(x_s) = \hat{V_1} \cdot \cos\left(\frac{x_s\pi}{\tau_p} - \alpha\right)$$

 α = 0: in U-axis



M.M.F. of three phase winding at "common mode current = zero"



a) Three phase sinusoidal currents

b) Three phase arbitrary currents



- In both cases a step like m.m.f. distribution with a dominant fundamental v = 1 is excited:
 - Double pole pitch $2\tau_p$ = wave length
 - Amplitude is derived from FOURIER-analysis: V_1

Fundamental moves with a) constant speed n_{syn} Fundamental amplitude is a) constant

```
b) arbitrary speed n(t) = \dot{\alpha}(t)/(2\pi \cdot p)
b) of arbitrary value \hat{V}_1(t)
```

- Here: ONLY THE FUNDAMENTAL ν = 1 IS FURTHER CONSIDERED !



"Arrow" for <u>fundamental</u> air gap magnetic flux density **B**

• Magnetic air gap field *B*: No slotting, no (or constant) iron saturation:



- **1.** Arrow length = \hat{B}
- 2. Arrow orientation = at maximum field
- 3. Arrow points in NORTH pole direction



$$B_1(x_s,t) = \mu_0 \cdot \frac{\hat{V}_1(t)}{\delta} \cdot \cos\left(\frac{x_s\pi}{\tau_p} - \alpha(t)\right) = \hat{B}(t) \cdot \cos\left(\frac{x_s\pi}{\tau_p} - \alpha(t)\right) = \hat{B}(t) \cdot \cos\left(\gamma_s - \alpha(t)\right)$$



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6. Space vector theory **Complex space vector = arrow of MMF and** *B***-fundamental**

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- Prof. Kovacs (Budapest): Complex co-ordinate system in machine cross section
- MMF fundamental may be represented as complex space phasor ("space vector"), lying in *machine cross section plane*.







Summary: M.M.F. space vector definition

- Three-phase system with arbitrary phase currents
- 2p-pole distributed winding leads to 2p-pole count also with arbitrary currents
- Movement of 2*p*-field follows the arbitrary time function of currents
- 2p-fundamental field described by space vector in magnitude and position
- Zero-sequence current system is here omitted,
- but if existing does not contribute to 2*p*-fundamental (see later in this chapter)





- 6. Space vector theory
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1st step: Single phase excitation: Magnetic alternating field



• Feeding the coil groups with sinusoidal alternating current i_c : Amplitude \hat{I}_c , frequency *f*, angular frequency $\omega = 2\pi f$, T = 1/f: period of oscillation

$$i_c(t) = \hat{I}_c \cdot \cos \omega t \implies B_\delta(x_s, t) = B_\delta(x_s) \cdot \cos \omega t$$

• Air gap field oscillates also sinusoidal with time, BUT maintains its spatial distribution (its shape = its distribution along x_s) ! The amplitude of (radial) field component $B_{\delta}(x_s, t)$ at locus x_s changes with time between positive and negative maximum value.





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FOURIER-series of field of one phase with pitched coil groups



$$\hat{V}_{\nu,ph}(t) = \frac{N}{2p} \cdot I(t) \cdot \frac{4}{\nu \pi} \cdot k_{p,\nu} \cdot k_{d,\nu}, \quad \nu = 1, 3, 5, \dots \quad N : \text{turns per phase}$$

- Phase current: Arbitrary current I(t), not: $I(t) = I_{rms} \cdot \sqrt{2} \cdot \cos(\omega t)$
- The MMF distribution $V_{\rm ph}(\gamma_s, t)$ (and hence the air gap field $B_{\delta}(\gamma_s, t)$) is a sum of pulsating, standing waves ("Pulsating field" with ~ I(t))

$$V_{ph}(\gamma_s, t) = \sum_{\nu=1,3,5,\dots}^{\infty} \hat{V}_{\nu,ph}(t) \cdot \cos(\nu \cdot \gamma_s)$$

- "Winding coefficient": $k_{w,v} = k_{p,v} \cdot k_{d,v}$
- Only fundamental v = 1 considered, at arbitrary current I(t): $\hat{V}_{1,ph}(t) = \frac{N}{2p} \cdot I(t) \cdot \frac{4}{\pi} \cdot k_{w,1}$
- At rated current $I_{\rm N}$:

$$\hat{V}_{1N,ph} = \frac{N}{2p} \cdot \hat{I}_N \cdot \frac{4}{\pi} \cdot k_{w,1}$$



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6. Space vector theory Calculation of m.m.f. space vector $\underline{V}(t)$ from arbitrary currents for a three-phase system



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2nd step: Fundamental of phase MMF distribution of phase U is directly proportional to the phase current value $I_{U}(t)$: $\hat{V}_{1U,ph}(t) = \hat{V}_{1N,ph} \cdot i_{U}(t)$

$$V_{1U}(x_s,t) = \hat{V}_{1U,ph}(t) \cdot \cos\left(\frac{x_s\pi}{\tau_p}\right) = \hat{V}_{1N,ph} \cdot \frac{I_U(t)}{\hat{I}_N} \cdot \cos\left(\frac{x_s\pi}{\tau_p}\right) = \hat{V}_{1N,ph} \cdot i_U(t) \cdot \cos(\gamma_s)$$



Calculation of m.m.f. space vectors for each phase





3rd step:

MMF fundamentals of phases V, W, excited by the arbitrary currents $I_v(t)$, $I_W(t)$, are spatially shifted by 120°, 240° (in el. degrees)!

$$V_{1V}(x_{s},t) = \hat{V}_{1N,ph} \cdot \frac{I_{V}(t)}{\hat{I}_{N}} \cdot \cos\left(\frac{x_{s}\pi}{\tau_{p}} - 2\pi/3\right) = \hat{V}_{1N,ph} \cdot i_{V}(t) \cdot \cos(\gamma_{s} - 2\pi/3) = V_{1V}(\gamma_{s},t)$$
$$V_{1W}(x_{s},t) = \hat{V}_{1N,ph} \cdot \frac{I_{W}(t)}{\hat{I}_{N}} \cdot \cos\left(\frac{x_{s}\pi}{\tau_{p}} - 4\pi/3\right) = \hat{V}_{1N,ph} \cdot i_{W}(t) \cdot \cos(\gamma_{s} - 4\pi/3) = V_{1W}(\gamma_{s},t)$$



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4th step:

Superposition of MMF fundamentals of all three phases U, V, W, and divide by 3/2 to get a per-unit value.



Calculation of MMF space vector $\underline{V}(t)$ from actual phase currents



$$\underline{a} = e^{j \cdot (2\pi/3)}$$

$$\underline{a}^{2} = e^{j \cdot (4\pi/3)}$$

$$\underbrace{a^{2} = e^{j \cdot (4\pi/3)}}_{i_{W}(t)}$$

$$\underbrace{v}_{i_{W}(t)}$$

$$\underbrace{\frac{2\pi}{3}}_{i_{U}(t)}$$

$$\underbrace{V_{1U}(\gamma_{s}, t) = \hat{V_{1N,ph}} \cdot i_{U}(t) \cdot \cos(-\gamma_{s} + 2\pi/3)}_{V_{1V}(\gamma_{s}, t) = \hat{V_{1N,ph}} \cdot i_{W}(t) \cdot \cos(-\gamma_{s} + 4\pi/3)}$$

$$\underbrace{\frac{a}{2} = e^{j \cdot (4\pi/3)}}_{i_{W}(t)}$$

$$\underbrace{V_{1}(\gamma_{s}, t) = \hat{V_{1N,ph}} \cdot \operatorname{Re}\left\{i_{U}(t) \cdot e^{-j\gamma_{s}} + i_{V}(t) \cdot e^{-j\gamma_{s}} \cdot \underline{a} + i_{W}(t) \cdot e^{-j\gamma_{s}} \cdot \underline{a}^{2}\right\}}$$

- Positions of the phasors of the MMF fundamentals of the three phases U, V, W : are spatially shifted by 120°, 240° (in el. degrees) \Rightarrow Multiplication with $\underline{a}, \underline{a}^2$ $\cos(-\gamma_s) = \operatorname{Re}\left\{e^{-j\gamma_s}\right\} \quad \cos(-\gamma_s + 2\pi/3) = \operatorname{Re}\left\{e^{-j\gamma_s} \cdot \underline{a}\right\} \quad \cos(-\gamma_s + 4\pi/3) = \operatorname{Re}\left\{e^{-j\gamma_s} \cdot \underline{a}^2\right\}$ $V_1(\gamma_s, t) = \hat{V}_{1N, ph} \cdot \operatorname{Re}\left\{i_U(t) + i_V(t) \cdot \underline{a} + i_W(t) \cdot \underline{a}^2\right] \cdot e^{-j\gamma_s}\right\} = \operatorname{Re}\left\{\underline{V}(t) \cdot e^{-j\gamma_s}\right\}$ $\underline{V}(t) = \hat{V}_{1N, ph} \cdot \left[i_U(t) + \underline{a} \cdot i_V(t) + \underline{a}^2 \cdot i_W(t)\right]$



6. Space vector theory Space fundamental MMF of a symmetrical 3-phase winding, fed by a 3-phase sinusoidal current system





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6. Space vector theory Per unit space vectors



• Fundamental wave m.m.f. amplitude at rated sinusoidal current system:

$$\hat{V}_{1,N} = \frac{3}{2} \cdot \hat{V}_{1N,ph} = \frac{\sqrt{2}}{\pi} \cdot \frac{3}{p} \cdot N \cdot k_{w,1} \cdot I_N$$
$$I_U(t) = I_N \cdot \sqrt{2} \cdot \cos(\omega t), I_V(t) = I_N \cdot \sqrt{2} \cdot \cos(\omega t - 2\pi/3), \ I_W(t) = I_N \cdot \sqrt{2} \cdot \cos(\omega t - 4\pi/3)$$

• M.M.F. space vector as <u>per unit value</u> of rated amplitude $\hat{V}_{1 N}$:

Arbitrary currents !

$$\underline{v}(t) = \frac{\underline{V}(t)}{\hat{V}_{1N}} = \frac{2}{3} \cdot \left[i_U(t) + \underline{a} \cdot i_V(t) + \underline{a}^2 \cdot i_W(t) \right]$$

- Current space vector: $\underline{I}(t) = \frac{2}{3} \cdot \left[I_U(t) + \underline{a} \cdot I_V(t) + \underline{a}^2 \cdot I_W(t) \right]$ (Definition !)
- The per unit current space vector is identical with per unit m.m.f. space vector:

$$\underline{i}(t) = \frac{\underline{I}(t)}{\widehat{I}_N} = \frac{2}{3} \cdot \left[i_U(t) + \underline{a} \cdot i_V(t) + \underline{a}^2 \cdot i_W(t) \right] = \underline{v}(t)$$



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Summary: M.M.F. space vector and phase currents

- Cross section plane of machine scaled with complex coordinate frame
- Space vector is a complex phasor in the cross-section coordinate frame
- Space vector formulation of 2*p*-fundamental field acc. to Prof. KOVACS
- Per-unit MMF space vector $\underline{v}(t)$ identical with per-unit current space vector $\underline{i}(t)$





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6. Space vector theory Phase shifter <u>a</u> : Calculation methods

0

Im

$$\underline{a}^{=e^{j\cdot(2\pi/3)}} \underbrace{a^{2} = e^{j\cdot(4\pi/3)}}_{1 = e^{j\cdot0}} \underbrace{e^{j} = e^{j\cdot(2\pi/3)}}_{1 = e^{j\cdot(2\pi/3)}} = e^{j\cdot(4\pi/3)} \underbrace{a^{3} = (e^{j\cdot(2\pi/3)})^{3} = e^{j\cdot(6\pi/3)} = e^{j\cdot2\pi} = 1}_{1/\underline{a} = (e^{j\cdot(2\pi/3)})^{-1} = e^{j\cdot(-2\pi/3)} = e^{j\cdot(4\pi/3)} = \underline{a}^{2}}_{1/\underline{a}^{2} = \underline{a} \Leftrightarrow \underline{a}^{3} = 1}_{1 + \underline{a} + \underline{a}^{2} = 0}$$

Example:

$$1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}} = 1 + \underline{a}^2 + \underline{a}^4 = 1 + \underline{a}^2 + \underline{a}^3 \cdot \underline{a} = 1 + \underline{a}^2 + \underline{a} = 0$$





Current space vector for three phase AC sine wave system



• Three phase sine wave current system:

a)
$$I_U(t) = \hat{I} \cdot \cos(\Omega \cdot t), I_V(t) = \hat{I} \cdot \cos(\Omega \cdot t - 2\pi/3), I_W(t) = \hat{I} \cdot \cos(\Omega \cdot t - 4\pi/3).$$

b) $I_U(t) = \hat{I} \cdot \frac{e^{j\Omega \cdot t} + e^{-j\Omega \cdot t}}{2}, I_V(t) = \hat{I} \cdot \frac{e^{j(\Omega \cdot t - 2\pi/3)} + e^{-j(\Omega \cdot t - 2\pi/3)}}{2}$
 $I_W(t) = \hat{I} \cdot \frac{e^{j(\Omega \cdot t - 4\pi/3)} + e^{-j(\Omega \cdot t - 4\pi/3)}}{2}$

• Current space vector:

$$\underline{I}(t) = \frac{2}{3} \cdot \left[\hat{I} \cdot \frac{e^{j\Omega \cdot t} + e^{-j\Omega \cdot t}}{2} + e^{j\frac{2\pi}{3}} \cdot \hat{I} \cdot \frac{e^{j(\Omega \cdot t - 2\pi/3)} + e^{-j(\Omega \cdot t - 2\pi/3)}}{2} + e^{j\frac{4\pi}{3}} \cdot \hat{I} \cdot \frac{e^{j(\Omega \cdot t - 4\pi/3)} + e^{-j(\Omega \cdot t - 4\pi/3)}}{2}}{2} \right] = \frac{2}{3} \cdot \left[\frac{3\hat{I}e^{j\Omega \cdot t}}{2} + \frac{\hat{I}e^{-j\Omega \cdot t}}{2} \cdot \left(\underbrace{1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} \right) \right] = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\Omega \cdot t}}_{0} \cdot \left(\underbrace{1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} \right) \right] = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\Omega \cdot t}}_{0} \cdot \left(\underbrace{1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} \right) \right] = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\Omega \cdot t}}_{0} \cdot \left(\underbrace{1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} \right) = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\Omega \cdot t}}_{0} \cdot \left(\underbrace{1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} \right) = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\Omega \cdot t}}_{0} \cdot \left(\underbrace{1 + e^{j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} \right) = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} = \underbrace{\hat{I} \cdot e^{j\frac{4\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\frac{4\pi}{3}} + e^{j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} + \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}_{0} = \underbrace{\hat{I}e^{-j\frac{8\pi}{3}}$$

• The current space vector of a three phase AC sine wave **system is a rotating vector of constant amplitude**, which is equal to the phase current amplitude. Rotation frequency is the electric frequency of the phase currents.



Current space vector at stationary operation





- Real fundamental field wave rotates in 2*p*-pole machine at $n_{syn} = f_s/p$
- Space vector rotates in 2p = 2-pole machine model with $n_{syn} = f_s$



Main flux linkage $\Psi_{h}(t)$ per phase (*m* = 3: U, V, W) DARMSTADT Main flux linkage per phase U: -W+U+V-U+W-V+U $\Psi_{h,U}(t) = \Psi_{h,UU}(t) + \Psi_{h,UV}(t) + \Psi_{h,UW}(t) =$ $au_{ m p}$ $= L_{h,ph} \cdot I_U(t) + M_{h,UV} \cdot I_V(t) + M_{h,UW} \cdot I_W(t)$ $\frac{M_{h,UV}}{L_{h,ph}} = \frac{M_{h,UV} \cdot i}{L_{h,ph} \cdot i} = \frac{-\Phi/2}{\Phi} = -\frac{1}{2}$ $\frac{M_{h,UW}}{L_{h,ph}} = -\frac{1}{2}$ \hat{B}_U $B_{\delta,\mathrm{U}}(\mathbf{x}) \sim I_{\mathrm{U}}$ ONLY fundamental field wave per phase considered \vec{x} $\Psi_{h,U}(t) = L_{h,ph} \cdot I_U - \frac{L_{h,ph}}{2} \cdot I_V - \frac{L_{h,ph}}{2} \cdot I_W$ Ф $B_{\delta,\mathrm{V}}(x) \sim I_{\mathrm{V}}$ $\hat{B}_V \sim I_V$ KIRCHHOFF's law: $I_{U}(t) + I_{V}(t) + I_{W}(t) = 0$ $-\Phi/2$ $I_U = -I_V - I_W$ $\Psi_{h,U} = L_{h,ph} \cdot I_U + \frac{L_{h,ph}}{2} \cdot (-I_V - I_W)$ $B_{\delta,W}(x) \sim I_W \qquad \hat{B}_W \sim I_W$ - *D*/2 $\Psi_{h,U}(t) = \frac{3L_{h,ph}}{2} \cdot I_U(t) = L_h \cdot I_U(t)$ \vec{x}

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Main flux linkage space vector $\underline{\Psi}_{h}(t)$







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Rotor current space vector



- Cage induction machine: Rotor contains Q_r rotor bars
- Slip-ring induction machine: Rotor contains the 3 rotor phase winding



	stator AC winding	rotor cage winding	Wound rotor
phase count	$m_{\rm s}=3$	$Q_{\rm r}$	$m_{\rm r}=3$
turns per phase	$N_{ m s}$	1/2	Nr
winding factor (v=1)	$k_{ m w,1,s}$	1	k _{w,1,r}

Magnetizing current space vector \underline{i}_{m}





- Stator current space vector \underline{i}_s is directly proportional to fundamental stator field wave
- Cage rotor is described by equivalent three-phase system
- Rotor current space vector $\underline{i'}_{r}$ represents rotor fundamental field wave
- Addition of stator and rotor space vector \underline{i}_{s} , $\underline{i'}_{r}$ in the same coordinate system \Rightarrow leads to magnetizing current space vector $\underline{i}_{m} \Rightarrow$ Gives resulting air-gap field wave, which is $\sim \underline{i}_{m} \sim \underline{\psi}_{h}$



6. Space vector theory Voltage space vector (1)



• Adding stator leakage flux linkage per phase yields stator flux linkage space vector: $\underline{\Psi}_{s}(t)$

e.g.:
$$\Psi_{sU}(t) = (L_h + L_\sigma) \cdot I_{sU}(t)$$
 per phase U, V, W
 $\underline{\Psi}_s(t) = (L_h + L_\sigma) \cdot \underline{I}_s(t)$

- At $R_s = 0$ the stator terminal voltage per phase is the resulting induced voltage per phase: $U_{sU}(t) = d \Psi_{sU}(t)/dt, U_{sV}(t) = d \Psi_{sV}(t)/dt, U_{sW}(t) = d \Psi_{sW}(t)/dt$ $\frac{d}{dt} \Psi_s(t) = \frac{2}{3} \cdot \frac{d}{dt} \left[\Psi_{sU}(t) + \underline{a} \cdot \Psi_{sV}(t) + \underline{a}^2 \cdot \Psi_{sW}(t) \right] = \frac{2}{3} \cdot \left[U_{sU}(t) + \underline{a} \cdot U_{sV}(t) + \underline{a}^2 \cdot U_{sW}(t) \right]$
- Definition of (stator) voltage space vector: $\underline{U}(t) = \frac{2}{3} \cdot \left[U_U(t) + \underline{a} \cdot U_V(t) + \underline{a}^2 \cdot U_W(t) \right]$
- Voltage space vector is an "artificial" quantity, which does not represent the real rotating field wave in the machine cross section. It is defined as analogous quantity!



Voltage space vector (2)









Example: Pulse width modulated inverter operation



Equivalent switching scheme of DC link voltage source inverter, connected to the two phases with switching transistor and free-wheeling diode

DC link voltage source inverter with switching transistors and free-wheeling diodes) e.g.: PM synchronous machine, R_s neglected: $U_d = u_{s,LL}(t) = d\psi_{s,LL} / dt$ $U_d - u_{p,LL} \approx L_{s,LL} \cdot di_s / dt$



Current ripple and chopped inverter voltage



Example: Block current operation



Block current commutation



Shaping of block current with hysteresis band control

For details see: Lecture "Motor development for electrical drive systems" Idealized shape of stator block current & synchronous back EMF u_p (= u_i) per phase





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Flux linkage space vector at inverter operation

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Example: Inverter-fed AC PM synchronous machine with **block-commutated** stator currents. Give stator flux linkage space vector WITHOUT rotor flux linkage space vector (= no rotor magnets)!









Summary: Current, voltage and flux linkage space vector

- Current space vector \underline{i} directly proportional to fundamental field amplitude \underline{B}
- Analogue definitions for voltage and flux linkage space vectors $\underline{u}, \, \underline{\psi}$
- Cage rotor described by equivalent three-phase system
- Rotor current space vector $\underline{i'_r}$ represents rotor field
- Addition of stator and rotor space vector \underline{i}_{s} , \underline{i}_{r} in the same coordinate system \Rightarrow leads to $\underline{i}_{m} \Rightarrow$ Gives resulting air-gap field $\underline{i}_{m} \sim \underline{\psi}_{h}$
- Orbit of space vector in cross-section plane is determined by phase current shape
- Sine wave current system: Space vector orbits on a circle





- 6. Space vector theory
 - 6.1 M.M.F. space vector definition
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 - 6.3 Current, flux linkage and voltage space vectors
 - 6.4 Space vector transformation
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 - 6.6 Magnetic energy



6. Space vector theory Space vector transformation



• Phase currents (voltages, flux linkages) U, V, W into α , β - components of space vector:





6. Space vector theory Space vector transformation (*Edith CLARKE*)



• Phase currents (voltages, flux linkages) of U, V, W into α , β - components of space vector: *CLARKE*'s transformation

$$\underline{I}(t) = \frac{2}{3} \cdot \left[I_U(t) + \underline{a} \cdot I_V(t) + \underline{a}^2 \cdot I_W(t) \right] = I_\alpha(t) + j \cdot I_\beta(t)$$
$$\underline{a} = -\frac{1}{2} + j \frac{\sqrt{3}}{2} = e^{j\frac{2\pi}{3}} \qquad \underline{a}^2 = -\frac{1}{2} - j \frac{\sqrt{3}}{2} = e^{j\frac{4\pi}{3}}$$



• α , β -components:

$$I_{\alpha}(t) = \frac{2}{3}I_{U}(t) - \frac{1}{3}(I_{V}(t) + I_{W}(t)) = I_{U}(t), * I_{\beta}(t) = (I_{V}(t) - I_{W}(t)) / \sqrt{3}$$

*): if $I_{\rm U} + I_{\rm V} + I_{\rm W} = 0$

- Three phase currents U, V, W a transformed into one space vector with two space vector components α , β with perpendicular directions ! = "Two-axes theory".



Two-axis theory corresponds to two-phase windings







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Space vector transformation for $I_0(t) = 0$ (1)



Clarke's matrix transformation:

$$\begin{aligned} \mathbf{U}, \mathbf{V}, \mathbf{W} \Rightarrow \alpha, \beta : \\ \begin{pmatrix} I_{\alpha}(t) \\ I_{\beta}(t) \end{pmatrix} &= \begin{pmatrix} \frac{2}{3} & -\frac{1}{3} & -\frac{1}{3} \\ 0 & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \end{pmatrix} \cdot \begin{pmatrix} I_{U}(t) \\ I_{V}(t) \\ I_{W}(t) \end{pmatrix} &= (A) \cdot \begin{pmatrix} I_{U}(t) \\ I_{W}(t) \\ I_{W}(t) \end{pmatrix} \end{aligned}$$
$$\begin{aligned} \alpha, \beta \Rightarrow \mathbf{U}, \mathbf{V}, \mathbf{W} : \\ (I_{U} + I_{V} + I_{W} = 0) \\ \begin{pmatrix} I_{U}(t) \\ I_{V}(t) \\ I_{W}(t) \end{pmatrix} &= \begin{pmatrix} 1 & 0 \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} \\ -\frac{1}{2} & -\frac{\sqrt{3}}{2} \end{pmatrix} \cdot \begin{pmatrix} I_{\alpha}(t) \\ I_{\beta}(t) \end{pmatrix} &= (A)^{-1} \cdot \begin{pmatrix} I_{\alpha}(t) \\ I_{\beta}(t) \\ I_{\beta}(t) \end{pmatrix} \end{aligned}$$
$$e.g.: I_{V} = -\frac{I_{\alpha}}{2} + \frac{\sqrt{3} \cdot I_{\beta}}{2} \end{aligned}$$



6. Space vector theory
Space vector transformation for
$$I_0(t) = 0$$
 (2)

$$\begin{bmatrix}
I_U \\
I_V \\
I_W
\end{bmatrix} = \begin{bmatrix}
1 & 0 \\
-\frac{1}{2} & \frac{\sqrt{3}}{2} \\
-\frac{1}{2} & -\frac{\sqrt{3}}{2} \\
-\frac{1}{2} & -\frac{\sqrt{3}}{2}
\end{bmatrix} \cdot \begin{bmatrix}
2 & -\frac{1}{3} & -\frac{1}{3} \\
0 & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \\
0 & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}}
\end{bmatrix} \cdot \begin{bmatrix}
I_U \\
I_V \\
I_W
\end{bmatrix} = \begin{bmatrix}
2 & -\frac{1}{3} & -\frac{1}{3} \\
-\frac{$$



Inverse space vector transformation: α , $\beta \Rightarrow U$, *V*, *W*



Valid for: $I_U(t) + I_V(t) + I_W(t) = 0$

$$I_U(t) = \operatorname{Re}\left\{\underline{I}(t)\right\}$$
$$I_V(t) = \operatorname{Re}\left\{\underline{a}^2 \cdot \underline{I}(t)\right\}$$
$$I_W(t) = \operatorname{Re}\left\{\underline{a} \cdot \underline{I}(t)\right\}$$

Proof:

$$I_{U}(t) = \operatorname{Re}\left\{\underline{I}(t)\right\} = \frac{2}{3} \cdot \left(I_{U}(t) + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{V}(t) + \operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{W}(t)\right) = \frac{2}{3} \cdot \left(I_{U}(t) - \frac{I_{V}(t)}{2} - \frac{I_{W}(t)}{2}\right) = \frac{2}{3} \cdot \left(I_{U}(t) + \frac{I_{U}(t)}{2}\right) = I_{U}(t)$$

$$I_{V} = \operatorname{Re}\left\{\underline{a}^{2}\underline{I}\right\} = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}^{3}\right\} \cdot I_{V} + \operatorname{Re}\left\{\underline{a}^{4}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{U} + \operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot I_{W}\right) = \frac{2}{3} \cdot \left(\operatorname{Re}\left\{\underline{a}^{2}\right\} \cdot \left$$

$$=\frac{2}{3} \cdot \left(-\frac{I_U + I_W}{2} + I_V\right) = \frac{2}{3} \cdot \left(\frac{I_V}{2} + I_V\right) = I_V$$

In the same way one proofs the last relationship $I_W(t) = \operatorname{Re}\{\underline{aI}(t)\}$.

Note: e.g.:
$$I_V = \operatorname{Re}\left\{\!\!\!\left[\frac{1}{2} \cdot \underline{I}\right]\!\!\!= \operatorname{Re}\left\{\!\!\left[\left(-\frac{1}{2} - j \cdot \frac{\sqrt{3}}{2}\right)\!\cdot \left(I_{\alpha} + j \cdot I_{\beta}\right)\!\!\right]\!\!\!= -\frac{I_{\alpha}}{2} + \frac{\sqrt{3} \cdot I_{\beta}}{2}\right]\!\!$$



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6. Space vector theory

Im

 $\frac{3}{2}$

Space vector transformation graphically

Example:

Star connected winding, per unit current values: $i_U = 0.3$, $i_V = 0.5$, $i_W = -0.8$

Im









Summary: Space vector transformation

- Coordinate system (Re, Im) of machine cross-section alternatively as α - β -system
- "Two axis"-theory in α - β -components
- Transformation from U, V, W into α - β -components = space vector transformation
- Edith CLARKE's transformation: Originally introduced for power systems,

not for electrical machines





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Influence of zero sequence current system



- In all three phases U, V, W the <u>same, identical</u> zero-sequence current $I_0(t)$ flows.
- Dynamic operation: Arbitrary time function: $I_0(t)$ is called COMMON MODE current !

Star connected winding	Delta connected winding	Star connected winding, neutral point connected
No common mode current	Common mode current flows <u>circulating</u> in delta connection, but is not visible in grid connections	Common mode current flows in each line and phase and with <u>3-times in neutral point</u> connection
$I_{II}(t) + I_{V}(t) + I_{W}(t) =$	Phase currents:	Phase currents:
$= 3 \cdot I_0(t) = 0$	$I_U(t) = I_V(t) = I_W(t) = I_0(t)$ Line currents:	$I_U(t) = I_V(t) = I_W(t) = I_0(t)$ Line currents:
$I_{0}(t) = 0$	$I_{L,UV}(t) = I_U(t) - I_V(t) = 0$ $I_{L,UV} \qquad I_0 \qquad I_V$	$I_{L,U}(t) = I_U(t) = I_0(t)$ Neutral current: $I_n(t) = I_U(t) + I_V(t) + I_W(t) =$ $= 3 \cdot I_0(t)$

Zero sequence (= common mode) current $I_0(t)$





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Result: One can always decompose a three-phase system

6. Space vector theory

Calculation of zero sequence current

Example:

- Star connected 3-phase winding, neutral point connected:
- At time *t*: Measured per unit currents: $i_U = 0.3$, $i_V = 0.5$, $i_W = -0.2$
- Per unit currents: $i(t) = I(t) / \hat{I}_N$
- Question: How big is zero sequence current a) in phases, b) in neutral clamp?

into a zero sequence system and a $_{n}I_{0}^{n}$ -free three phase system.

a)
$$i_0(t) = \frac{1}{3} \cdot (i_U + i_V + i_W) = \frac{1}{3} \cdot (0.3 + 0.5 - 0.2) = \underline{0.2}$$

b) $i_n(t) = 3i_0 = 3 \cdot 0.2 = 0.6$

Result: In neutral clamp flows 60% of rated current !

The "common-mode free" current values are:

 $i_{US} = i_U - i_0 = 0.3 - 0.2 = 0.1$ $i_{VS} = i_V - i_0 = 0.5 - 0.2 = 0.3$

$$i_{WS} = i_W - i_0 = -0.2 - 0.2 = -0.4$$



Air gap magnetic field due to zero sequence current



• **Example:** $i_U = 0.3, i_V = 0.5, i_W = -0.2 \implies i_0(t) = \frac{1}{3} \cdot (0.3 + 0.5 - 0.2) = \underline{0.2}$

Magnetic air gap field, excited by zero sequence current (= flowing in all three phases). <u>Here:</u> Single layer winding with q = 2 slots per pole and phase:



• Ampere's law yields a zero-sequence M.M.F. distribution $V_0(x_s)$ with three pole pairs instead of one along double pole pitch $2\tau_p$.



Effect of zero sequence air gap field



- Zero sequence field is not travelling, but STANDING, but time varying acc. to $i_0(t)$.
- 3 pole pairs of stator with one pole pair of rotor (e.g. PM machine) do not generate torque.
- Zero sequence current excites flux waves, which <u>do not contribute</u> to torque generation of fundamental wave.
- Space vector of zero sequence current is therefore zero.

• Proof:
$$I_0(t) = \frac{2}{3} \cdot \left(I_0(t) + \underline{a} \cdot I_0(t) + \underline{a}^2 \cdot I_0(t) \right) = \frac{2}{3} \cdot I_0(t) \cdot \left(1 + \underline{a} + \underline{a}^2 \right) =$$

$$= \frac{2}{3} \cdot I_0(t) \cdot \left(1 + e^{j\frac{2\pi}{3}} + e^{j\frac{4\pi}{3}} \right) = 0$$

 <u>Facit</u>: Additional losses, pulsating radial forces (leading to noise or vibration), in cage induction machines also parasitic <u>braking</u> torque may occur, so AVOID common mode current !





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Effect of zero sequence current in 3-leg transformers

Symmetrical current system e.g. Yy0



Common-mode current system possible e.g. at Dy



Per phase at primary side at no-load ($I_2 = 0$):

$$\Psi_{1,ph}(t) = (L_{1\sigma} + L_{1h}) \cdot I_{1,ph}(t)$$

Big magnetizing inductance $L_{1h} \sim \mu_{Fe}$! (No air-gap, moderate iron saturation)

 $\Psi_{1,ph,0}(t) = L_{1,0} \cdot I_{1,ph,0}(t)$

Small zero-sequence inductance $L_{1,0} \sim \mu_0$! (Flux lines pass via air = similar to stray flux!)







Effect of zero sequence current in 5-leg transformers

Common-mode current system



$$L_{1,0} = l_V = l_W = l_0$$

$$L_{1,0} < L_{1h}$$

Per phase at primary side at no-load ($I_2 = 0$): $\Psi_{1,ph,0}(t) = L_{1,0} \cdot I_{1,ph,0}(t)$

Big zero-sequence inductance $L_{1,0} \sim \mu_{\text{Fe}}$! Flux lines pass partially via the outer legs and not via air \Rightarrow zero-sequence flux bigger than stray flux!

- (i) In 5-leg transformers the zero-sequence inductance $L_{1,0} < L_{1h}$ is much bigger than in 3-leg transformers
- (ii) and decreases with increasing current $I_0(t)$ due to iron saturation of the outer legs





Space vector CLARKE transformation including $I_0(t)$

Clarke's matrix transformation from 3-phase system U, V, W to two-axes system α , β and zero-sequence system 0:

$$\mathbf{U}, \mathbf{V}, \mathbf{W} \Rightarrow \alpha, \beta, \mathbf{0}: \qquad \begin{pmatrix} I_{\alpha}(t) \\ I_{\beta}(t) \\ I_{0}(t) \end{pmatrix} = \begin{pmatrix} \frac{-}{3} & -\frac{-}{3} & -\frac{-}{3} \\ 0 & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \\ \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} \\ \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \end{pmatrix} \cdot \begin{pmatrix} I_{U}(t) \\ I_{V}(t) \\ I_{W}(t) \end{pmatrix} = (A) \cdot \begin{pmatrix} I_{U}(t) \\ I_{W}(t) \end{pmatrix}$$

$$\alpha, \beta, \mathbf{0} \Rightarrow \mathbf{U}, \mathbf{V}, \mathbf{W}: \qquad \begin{pmatrix} I_{U}(t) \\ I_{V}(t) \\ I_{W}(t) \end{pmatrix} = \begin{pmatrix} 1 & 0 & 1 \\ -\frac{1}{2} & \frac{\sqrt{3}}{2} & 1 \\ -\frac{1}{2} & -\frac{\sqrt{3}}{2} & 1 \end{pmatrix} \cdot \begin{pmatrix} I_{\alpha}(t) \\ I_{\beta}(t) \\ I_{0}(t) \end{pmatrix} = (A)^{-1} \cdot \begin{pmatrix} I_{\alpha}(t) \\ I_{\beta}(t) \\ I_{0}(t) \end{pmatrix}$$

<u>Note:</u> Transformation from *m*-phase system to two-axes system α , β and zero-sequence system 0 yields a (3 x *m*)-matrix!



6. Space vector theory *CLARKE* transformation identities







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Summary: Influence of zero sequence current system

- Zero-sequence current system excites a 6p-air gap field
- Hence zero sequence system does not contribute to space vectors, as these are related to 2*p*-air gap field waves
- Insulated star-point: No zero sequence current system possible
- (3x3)-CLARKE's transformation includes zero sequence system
- Zero sequence systems should be avoided: otherwise additional losses, pulsating torque, vibration forces, noise, …





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 - (6.6 Magnetic energy: leads via power balance to torque equation)

